

Abstract interpolation problem and four–block problem

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We treat a four-block problem as an abstract interpolation problem with an appropriately chosen data. This approach enables us to obtain a description of all solutions for the four-block problem.

In this note we suggest a way of the inclusion of the four-block problem [6, Ch. IX, Sect. 4] into the framework of the abstract interpolation problem [8–10]. We would like to note that one can investigate the four-block problem with the help of the commutant lifting theorem [6, 13, Ch. II, Sect. 2]. This approach is presented in the monograph [6] and in the papers [5, 7].

We should mention here that the relation between general commutant lifting and a coordinate-free version of the abstract interpolation problem is discussed in [3].

The reasoning we are going to present is quite similar to that of [3, 11].

1. Abstract interpolation problem

1.1. We would like to begin the exposition with a brief recalling of the concept of the abstract interpolation problem (or the (AI)-problem, in abbreviated form). The formulation of the problem and its detailed discussion can be found in [8–10]. The second paper contains also the applications of the problem to the generalized bitangential Schur–Nevanlinna–Pick problem and j -inner-outer factorization as well as a wide bibliography on the question. We would like to mention the paper [4] that deals with a very close circle of ideas.

Let us have a complex linear space X and a nonnegative quadratic form D on the space, $D(x, x) \geq 0, \forall x \in X$. Let us assume that the identity holds

$$D(T_2x, T_2y) - D(T_1x, T_1y) = (M_1x, M_1y)_{\varepsilon_1} - (M_2x, M_2y)_{\varepsilon_2}, \forall x, y \in X. \quad (1.1.1)$$

Here T_1 and T_2 are some operators, defined on X , \mathcal{E}_1 and \mathcal{E}_2 are separable Hilbert spaces, M_1 and M_2 are linear maps, acting from X to \mathcal{E}_1 and from X to \mathcal{E}_2 , respectively. For the sake of brevity we call the set $\{X, \mathcal{E}_1, \mathcal{E}_2; D, T_1, T_2, M_1, M_2\}$ an interpolation data of the (AI)-problem.

We should give some definitions to introduce the notion a solution of the problem. Let $\omega(\zeta)$ be a $[\mathcal{E}_1, \mathcal{E}_2]$ -valued analytic contractive function on the unit disk $\mathbf{D} = \{\zeta \in \mathbf{C} : |\zeta| < 1\}$. We consider an operator-valued function

$$\begin{bmatrix} 1_{\mathcal{E}_1} & \omega^* \\ \omega & 1_{\mathcal{E}_2} \end{bmatrix} : \begin{bmatrix} L^2(\mathcal{E}_1) \\ L^2(\mathcal{E}_2) \end{bmatrix} \rightarrow \begin{bmatrix} L^2(\mathcal{E}_1) \\ L^2(\mathcal{E}_2) \end{bmatrix}.$$

It is almost obvious that the defined map is nonnegative, we denote its square root by $\Sigma_\omega, \Sigma_\omega = \begin{bmatrix} 1_{\mathcal{E}_1} & \omega^* \\ \omega & 1_{\mathcal{E}_2} \end{bmatrix}^{1/2}$. Next, we define a subspace L_ω^2 of $L^2(\mathcal{E}_1 \oplus \mathcal{E}_2)$ by the equality

$$L_\omega^2 = \overline{\Sigma_\omega \begin{bmatrix} L^2(\mathcal{E}_1) \\ L^2(\mathcal{E}_2) \end{bmatrix}}.$$

We set

$$H_\omega = L_\omega^2 \ominus \left\{ \Sigma_\omega \begin{bmatrix} H^2(\mathcal{E}_1) \\ 0 \end{bmatrix} \oplus \Sigma_\omega \begin{bmatrix} 0 \\ H_-^2(\mathcal{E}_2) \end{bmatrix} \right\} \quad (1.1.2)$$

now. Here $H_-^2(\mathcal{E}_2)$ stands for $L^2(\mathcal{E}_2) \ominus H^2(\mathcal{E}_2)$, $H^2(\mathcal{E}_2)$ is the standard Hardy space of vector-valued functions on the unit disk \mathbf{D} . The space H_ω is called de Branges–Rovnyak space, associated with function ω .

Let us observe that a vector $f = \begin{bmatrix} f_1 \\ f_2 \end{bmatrix}$ from L_ω^2 lies in H_ω if and only if the vector $\begin{bmatrix} f_- \\ f_+ \end{bmatrix} = \Sigma_\omega f$ is in $\begin{bmatrix} H_-^2(\mathcal{E}_1) \\ H^2(\mathcal{E}_2) \end{bmatrix}$. The assertion can be easily deduced from the equality

$$0 = \left(\begin{bmatrix} f_1 \\ f_2 \end{bmatrix}, \Sigma_\omega \begin{bmatrix} h_+ \\ h_- \end{bmatrix} \right) = \left(\Sigma_\omega \begin{bmatrix} f_1 \\ f_2 \end{bmatrix}, \begin{bmatrix} h_+ \\ h_- \end{bmatrix} \right),$$

where $h_+ \in H^2(\mathcal{E}_1), h_- \in H^2(\mathcal{E}_2)$.

By a solution of the (AI)-problem we mean a pair (ω, F) , where ω is a $[\mathcal{E}_1, \mathcal{E}_2]$ -valued contractive analytic on \mathbf{D} function and F is a linear map from X to H_ω , possessing the properties

$$1) \quad \bar{t} \left(FT_1 x \oplus \Sigma_\omega \begin{bmatrix} M_1 x \\ 0 \end{bmatrix} \right) = FT_2 x \oplus \Sigma_\omega \begin{bmatrix} 0 \\ \bar{t} M_2 x \end{bmatrix}, \quad \forall x \in X, \quad (1.1.3)$$

$$2) \quad \|Fx\|_2^2 \leq D(x, x), \quad \forall x \in X. \quad (1.1.4)$$

The map F is called the Fourier representation of the (AI)-problem. We would like to comment on relation (1.1.3). It can be readily seen that the summands

in both sides of the equality are really orthogonal one to another. Indeed, FT_1x and FT_2x lie in H_ω while $\Sigma_\omega \begin{bmatrix} M_1x \\ 0 \end{bmatrix} \in \Sigma_\omega \begin{bmatrix} H^2(\mathcal{E}_1) \\ 0 \end{bmatrix}$ and $\Sigma_\omega \begin{bmatrix} 0 \\ \bar{t}M_2x \end{bmatrix} \in \Sigma_\omega \begin{bmatrix} 0 \\ H^2_-(\mathcal{E}_2) \end{bmatrix}$, the rest is clear from the definition of H_ω . Further, we rewrite equality (1.1.3) in the form

$$\bar{t} \left(\begin{bmatrix} F_- \\ F_+ \end{bmatrix} T_1x \oplus \begin{bmatrix} 1_{\mathcal{E}_1} & \omega^* \\ \omega & 1_{\mathcal{E}_2} \end{bmatrix} \begin{bmatrix} M_1x \\ 0 \end{bmatrix} \right) = \begin{bmatrix} F_- \\ F_+ \end{bmatrix} T_2x \oplus \begin{bmatrix} 1_{\mathcal{E}_1} & \omega^* \\ \omega & 1_{\mathcal{E}_2} \end{bmatrix} \begin{bmatrix} 0 \\ \bar{t}M_2x \end{bmatrix}, \quad (1.1.5)$$

or, equivalently,

$$\begin{bmatrix} F_- \\ F_+ \end{bmatrix} T_1x = t \begin{bmatrix} F_- \\ F_+ \end{bmatrix} T_2x - \begin{bmatrix} 1_{\mathcal{E}_1} & \omega^* \\ \omega & 1_{\mathcal{E}_2} \end{bmatrix} \begin{bmatrix} M_1x \\ -M_2x \end{bmatrix}, \quad \forall x \in X. \quad (1.1.6)$$

Here $\begin{bmatrix} F_- \\ F_+ \end{bmatrix} x$ stands for $\Sigma_\omega Fx, x \in X$, and $t \in \mathbf{T} = \{\zeta \in \mathbf{C} : |\zeta| = 1\}$ stands for the independent variable on the unit circle.

1.2. We need more details about the inner machinery of the (AI)-problem. Firstly, it is possible to build a new Hilbert space that preserves, in essential, the structure of the space X . We define an antilinear functional \bar{x} with the help of an arbitrary vector x from X by formula

$$\bar{x}(y) = D(x, y), \quad \forall y \in X.$$

Denote the linear space of the functionals by \bar{X} . The space \bar{X} can be equipped with a scalar product in the following way:

$$(\bar{x}, \bar{y})_D = D(x, y), \quad \forall \bar{x}, \bar{y} \in \bar{X}.$$

The scalar product is well-defined because of the nonnegativity of D . Now we should only take a closure of \bar{X} in the metric to get a Hilbert space, $H = \text{clos}_{(\cdot, \cdot)_D} \bar{X}$.

We are going to define a certain isometry, acting on the Hilbert space H . The solutions of the (AI)-problem will be characterized in terms of unitary extensions of this isometry. Let us put

$$d_V = \{\overline{T_1x} \oplus M_1x : x \in X\} \subset H \oplus \mathcal{E}_1,$$

$$\Delta_V = \{\overline{T_2x} \oplus M_2x : x \in X\} \subset H \oplus \mathcal{E}_2.$$

Identity (1.1.1) states that the map, defined by the formula

$$V : d_V \rightarrow \Delta_V, \quad V(\overline{T_1x} \oplus M_1x) = \overline{T_2x} \oplus M_2x, \quad x \in X, \quad (1.2.1)$$

is an isometry. We refer to this operator as to the (AI)-isometry, also it is convenient to denote its defect subspaces by

$$N_{d_V} = (H \oplus \mathcal{E}_1) \ominus d_V \text{ and } N_{\Delta_V} = (H \oplus \mathcal{E}_2) \ominus \Delta_V. \quad (1.2.2)$$

1.3. In this subsection we demonstrate how a unitary extension of the (AI)-isometry generates a solution of (AI)-problem .

Let K be a Hilbert space that contains H as its subspace. By a unitary extension of an isometry $V, V : H \oplus \mathcal{E}_1 \rightarrow H \oplus \mathcal{E}_2$, we mean a unitary operator U acting from $K \oplus \mathcal{E}_1$ to $K \oplus \mathcal{E}_2$ in such a way that $U|_{d_V} = V$.

We prefer to treat the unitary operator as a unitary colligation with coefficient spaces $\mathcal{E}_1, \mathcal{E}_2$ and the state space K . In other words, we decompose the operator U into the blocks

$$U = \begin{bmatrix} A & B \\ C & D \end{bmatrix} : \begin{bmatrix} K \\ \mathcal{E}_1 \end{bmatrix} \rightarrow \begin{bmatrix} K \\ \mathcal{E}_2 \end{bmatrix}.$$

Following [1], we define the characteristic function of the colligation by the formula

$$\omega(\zeta) = D + \zeta C(I - \zeta A)^{-1} B. \quad (1.3.1)$$

Evidently, the function is analytic on the unit disk \mathbf{D} . Moreover, one can prove that the function is contractive [1, 2, 10]. Further, we define a map \mathbf{F} with the help of the relations

$$\mathbf{F}(\zeta) = \begin{bmatrix} \mathbf{F}_-(\zeta) \\ \mathbf{F}_+(\zeta) \end{bmatrix} = \begin{bmatrix} \bar{\zeta} B^* (I - \bar{\zeta} A^*)^{-1} \\ C(I - \zeta A)^{-1} \end{bmatrix} : K \rightarrow \begin{bmatrix} \mathcal{E}_1 \\ \mathcal{E}_2 \end{bmatrix}. \quad (1.3.2)$$

The map is called a Fourier representation of the state space K . The following proposition describes the main properties of the representation.

Proposition 1.1. (*[2, 8–10]*) *A Fourier representation, defined by (1.3.2), has the properties:*

1) *for every k from K the vector $\mathbf{F}k$ lies in $\Sigma_\omega H_\omega$, so the map $\Sigma_\omega^{(-1)} \mathbf{F}$, acting from K to H_ω , is well-defined.*

2)

$$\bar{t} \left(\mathbf{F} P_K \oplus \begin{bmatrix} 1_{\mathcal{E}_1} & \omega^* \\ \omega & 1_{\mathcal{E}_2} \end{bmatrix} \begin{bmatrix} 1_{\mathcal{E}_1} \\ 0 \end{bmatrix} P_{\mathcal{E}_1} \right) = \mathbf{F} P_K U \oplus \begin{bmatrix} 1_{\mathcal{E}_1} & \omega^* \\ \omega & 1_{\mathcal{E}_2} \end{bmatrix} \begin{bmatrix} 0 \\ 1_{\mathcal{E}_2} \end{bmatrix} P_{\mathcal{E}_2} U.$$

3)

$$\|\Sigma_\omega^{(-1)} \mathbf{F}k\|_2^2 \leq \|k\|_K^2, \forall k \in K.$$

It seems to be relevant to compare 2) with (1.1.5) and (1.2.1).

Assign $Fx = \Sigma_\omega^{(-1)} \mathbf{F}\bar{x}, x \in X$. Since properties 2) and 3) of the previous proposition turn into (1.1.3) and (1.1.4), we deduce that the pair (ω, F) is a solution of the (AI)-problem.

The cornerstone of the approach is that, conversely, every solution (ω, F) of the (AI)-problem arises in this way (see [8, 10]).

1.4. Let \mathcal{N}_1 and \mathcal{N}_2 be copies of the defect subspaces of the (AI)-isometry. This time we obtain a unitary extension of the isometry by enlarging the coefficient spaces instead of the state space. Namely, we consider a unitary operator U_0 mapping $H \oplus (\mathcal{E}_1 \oplus \mathcal{N}_2)$ onto $H \oplus (\mathcal{E}_2 \oplus \mathcal{N}_1)$. Note that

$$H \oplus (\mathcal{E}_1 \oplus \mathcal{N}_2) = (H \oplus \mathcal{E}_1) \oplus \mathcal{N}_2 = d_V \oplus N_{d_V} \oplus \mathcal{N}_2,$$

$$H \oplus (\mathcal{E}_2 \oplus \mathcal{N}_1) = (H \oplus \mathcal{E}_2) \oplus \mathcal{N}_1 = \Delta_V \oplus N_{\Delta_V} \oplus \mathcal{N}_1.$$

This remark permits to introduce the universal extension of the (AI)-isometry (see [1, 8, 9]) as

$$U_0|_{d_V} = V, \quad U_0|_{N_{d_V}} = id : N_{d_V} \rightarrow \mathcal{N}_1, \quad U_0|_{\mathcal{N}_2} = id : \mathcal{N}_2 \rightarrow N_{\Delta_V}.$$

Let us consider the characteristic function and the Fourier representation of the unitary colligation

$$S(\zeta) = D_0 + \zeta C_0(I - \zeta A_0)^{-1} B_0,$$

$$G(\zeta) = \begin{bmatrix} G_-(\zeta) \\ G_+(\zeta) \end{bmatrix} = \begin{bmatrix} \bar{\zeta} B_0^*(I - \bar{\zeta} A_0^*)^{-1} \\ C_0(I - \zeta A_0)^{-1} \end{bmatrix} : H \rightarrow \Sigma_S H_S,$$

where

$$U_0 = \begin{bmatrix} A_0 & B_0 \\ C_0 & D_0 \end{bmatrix} : \begin{bmatrix} H \\ \mathcal{E}_1 \oplus \mathcal{N}_2 \end{bmatrix} \rightarrow \begin{bmatrix} H \\ \mathcal{E}_2 \oplus \mathcal{N}_1 \end{bmatrix}.$$

It is convenient to introduce the block decompositions of the just defined maps:

$$G_+ = \begin{bmatrix} G_{+2} \\ G_{+1} \end{bmatrix} : H \rightarrow \begin{bmatrix} \mathcal{E}_2 \\ \mathcal{N}_1 \end{bmatrix}, \quad G_- = \begin{bmatrix} G_{-1} \\ G_{-2} \end{bmatrix} : H \rightarrow \begin{bmatrix} \mathcal{E}_1 \\ \mathcal{N}_2 \end{bmatrix},$$

$$S = \begin{bmatrix} s_0 & s_2 \\ s_1 & s \end{bmatrix} : \begin{bmatrix} \mathcal{E}_1 \\ \mathcal{N}_2 \end{bmatrix} \rightarrow \begin{bmatrix} \mathcal{E}_2 \\ \mathcal{N}_1 \end{bmatrix}.$$

Before stating the theorem, let us denote by $\mathbf{B}[\mathcal{N}_1, \mathcal{N}_2]$ the set of all $[\mathcal{N}_1, \mathcal{N}_2]$ -valued analytic contractive functions on the unit disk \mathbf{D} .

Theorem 1.1. ([8–10]) *All solutions of the (AI)-problem are parametrized by functions ε from $\mathbf{B}[\mathcal{N}_1, \mathcal{N}_2]$. The correspondence is given by formulas:*

$$\omega = s_0 + s_2(1_{\mathcal{N}_2} - \varepsilon s)^{-1} \varepsilon s_1, \tag{1.4.1}$$

$$\begin{bmatrix} F_-(\zeta)x \\ F_+(\zeta)x \end{bmatrix} = \begin{bmatrix} G_{-1} + s_1^*(1_{\mathcal{N}_1} - \varepsilon^* s^*)^{-1} \varepsilon^* G_{-2} \\ G_{+2} + s_2(1_{\mathcal{N}_2} - \varepsilon s)^{-1} \varepsilon G_{+1} \end{bmatrix} : H \rightarrow \Sigma_\omega H_\omega. \tag{1.4.2}$$

2. Formulation of the four-block problem

2.1. Let E_1, E'_1, E_2, E'_2, F be separable Hilbert spaces. Let

$$\Omega_0 = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \tag{2.1.1}$$

be a fixed function from $L^\infty(E_1 \oplus E_2, E'_1 \oplus E'_2)$. The operator-valued function maps $\begin{bmatrix} L^2(E_1) \\ L^2(E_2) \end{bmatrix}$ into $\begin{bmatrix} L^2(E'_1) \\ L^2(E'_2) \end{bmatrix}$ via multiplication. Let a $[F, E'_1]$ -valued inner function be given, $\theta \in H^\infty(F, E'_1)$. By the four-block problem [6, Ch. IX, Sect. 4] one understands the problem of finding δ ,

$$\begin{aligned} \delta &= \inf \left\| \begin{bmatrix} A + \theta H^\infty(E_1, F) & B \\ C & D \end{bmatrix} \right\|_\infty \\ &= \inf \left\{ \left\| \begin{bmatrix} A + \theta \Psi & B \\ C & D \end{bmatrix} \right\|_\infty : \Psi \in H^\infty(E_1, F) \right\}, \end{aligned}$$

and characterizing all optimal Ψ_* from $H^\infty(E_1, F)$.

A detailed discussion of the problem from the commutant lifting theorem point of view can be found in the already mentioned monograph [6, Ch. IX, Sect. 4]. The same source contains applications of the problem to the robust control theory as well as the history of the question.

In particular, the construction permits to calculate the value δ as a norm of a certain Hankel type operator. We need some definitions to formulate the result.

We define an inner function by the formula

$$\Phi = \begin{bmatrix} \theta \\ 0 \end{bmatrix} : L^2(F) \rightarrow \begin{bmatrix} L^2(E'_1) \\ L^2(E'_2) \end{bmatrix}, \Phi^* \Phi = 1_F.$$

Then we consider spaces

$$H = \begin{bmatrix} H^2(E_1) \\ L^2(E_2) \end{bmatrix} \text{ and } H' = \begin{bmatrix} L^2(E'_1) \\ L^2(E'_2) \end{bmatrix} \ominus \Phi H^2(F). \tag{2.1.2}$$

We denote the operator of the orthogonal projection from $L^2(E'_1) \oplus L^2(E'_2)$ to H' by $P_{H'}$. Further we set operators T and T' to be

$$Th = th, h \in H, \text{ and } T'h' = P_{H'}th', h' \in H'.$$

Let us define an operator Γ , mapping H to H' by the formula

$$\Gamma h = P_{H'}\Omega_0 h, h \in H. \tag{2.1.3}$$

Clearly, the operator is in the commutant of T and T' . Indeed,

$$\Gamma Th = \Gamma th = P_{H'}\Omega_0 th = P_{H'}t\Omega_0 h = P_{H'}tP_{H'}\Omega_0 h = T'\Gamma h, \forall h \in H.$$

Theorem 2.2. ([6, Ch. IX, Sect. 4]) Let Ω_0 be an operator-valued function defined by (2.1.1). Then

$$\inf \|\Omega_0 + \begin{bmatrix} \theta H^\infty(E_1, F) & 0 \\ 0 & 0 \end{bmatrix}\|_\infty = \|\Gamma\|.$$

2.2. We would like to study a question that slightly differs from the problem discussed in the previous subsection. Assuming that $\|\Gamma\| \leq 1$, we are interested in describing all operator-valued functions Ω from $L^\infty(E_1 \oplus E_2, E'_1 \oplus E'_2)$ which define the same operator Γ via (2.1.3) while their norms satisfy the inequality $\|\Omega\|_\infty \leq 1$.

More precisely, let an operator Γ , acting from H to H' , be given. We suppose that

$$\Gamma T = T' \Gamma \text{ on } H. \tag{2.2.1}$$

The problem is to parametrize the set of operator-valued functions $\Omega \in L^\infty(E_1 \oplus E_2, E'_1 \oplus E'_2)$ such that

$$1) \quad \Gamma h = P_{H'} \Omega h, \forall h \in H, \tag{2.2.2}$$

$$2) \quad \|\Omega\|_\infty \leq 1. \tag{2.2.3}$$

Hereafter we will refer to the problem as to the four-block problem (or, to be brief, to the (4B)-problem). The function Ω possessing properties (2.2.2) and (2.2.3), will be called a solution of the (4B)-problem.

2.3. The purpose of this subsection is to put the (4B)-problem into the framework of the scheme of the abstract interpolation.

It can be easily done by exploiting commutativity relation (2.2.1). We have

$$\Gamma T h = T' \Gamma h = P_{H'} t \Gamma h = t \Gamma h - \Phi M_2 h \tag{2.3.1}$$

by definition of $P_{H'}$ (see Subsect. 2.1). Here M_2 is an operator from H to $H^2(F)$. We rewrite (2.3.1) in the equivalent form

$$t \Gamma h = \Gamma t h + \Phi M_2 h$$

and we note that $\Gamma T h \perp \Phi M_2 h$, because $\Gamma T h \in H'$ and $\Phi M_2 h \perp H'$. This remark immediately implies that

$$(t \Gamma h, \Phi g) = (\Phi M_2 h, \Phi g) = (M_2 h, g) \quad \forall h \in H, \forall g \in H^2(F). \tag{2.3.2}$$

In particular, the equality shows that $(\Phi^* t \Gamma h - M_2 h, g) = 0 \quad \forall g \in H^2(F)$, or, consequently, $M_2 h = P_+ t \Phi^* \Gamma h \in F$. Furthermore, this enables us to calculate the expression in the following way:

$$\begin{aligned} ((I - \Gamma^* \Gamma) h_1, h_2) &= ((I - \Gamma^* \Gamma) t h_1, t h_2) = (\Gamma t h_1, \Gamma t h_2) - (\Gamma h_1, \Gamma h_2) \\ &= (t \Gamma h_1, t \Gamma h_2) - (t \Gamma h_1, \Phi M_2 h_2) - (\Phi M_2 h_1, t \Gamma h_2) \\ &+ (\Phi M_2 h_1, \Phi M_2 h_2) - (\Gamma h_1, \Gamma h_2) = -(M_2 h_1, M_2 h_2). \end{aligned}$$

We have applied relation (2.3.2) with $g = M_2h_2$ and its conjugate with $g = M_2h_1$ to get the latter equality.

Let us assign

$$\begin{aligned} D &= I - \Gamma^*\Gamma \geq 0 \text{ (because } \|\Gamma\| \leq 1), \\ M_2h &= P_+t\Phi^*\Gamma h : H \rightarrow F, \quad M_1h = 0, \\ T_1 &= T = t \cdot _, \quad T_2 = id \text{ on } H, \\ X &= H, \quad \mathcal{E}_1 = \{0\}, \quad \mathcal{E}_2 = F. \end{aligned} \tag{2.3.3}$$

Hence, the new setting transforms the previous equality to identity (1.1.1). It means that we have reduced the given (4B)-problem to an (AI)-problem with appropriate data. We will call the problem the (AI')-problem to stress the particular choice of interpolation data (2.3.3).

We want to finish this subsection with an explicit formula for the operator of the orthogonal projection $P_{H'}$. Arguing as during the computation of M_2h , one can get that

$$P_{H'}g = g - \Phi P_+ \Phi^*g, \quad g \in L^2(E'_1) \oplus L^2(E'_2), \tag{2.3.4}$$

where P_+ is the Riesz projection acting from $L^2(F)$ to $H^2(F)$ (see also [12, Ch. 2] or [14]). The relation will be useful in the next section.

3. Correspondence between the solutions of the (4B)-problem and the (AI')-problem

3.1. We begin this subsection with the following observation. Let (ω, F) be an arbitrary solution of the (AI')-problem. Since $\mathcal{E}_1 = \{0\}$ and ω is a $[\mathcal{E}_1, \mathcal{E}_2]$ -valued function, we automatically obtain that $\omega \equiv 0$ on the unit disk \mathbf{D} . Then, by definition (1.1.2)

$$H_\omega = \left[\begin{array}{c} \{0\} \\ L^2(F) \end{array} \right] \ominus \left[\begin{array}{c} \{0\} \\ H_-^2(F) \end{array} \right] = \left[\begin{array}{c} \{0\} \\ H^2(F) \end{array} \right] \simeq H^2(F).$$

Consequently, the first component F_- of the Fourier representation is identically equal to zero, while the second one maps H into $H^2(F)$ and satisfies the relations

$$1) \quad F_+th = tF_+h + M_2h, \quad \forall h \in H, \tag{3.1.1}$$

$$2) \quad (F_+h, F_+h) \leq ((I - \Gamma^*\Gamma)h, h), \quad \forall h \in H. \tag{3.1.2}$$

To derive them, we substitute the explicit expressions from (2.3.3) into formulas (1.1.4) and (1.1.6).

3.2. The connection between solutions of the (AI')-problem and the (4B)-problem is given by the following proposition.

Proposition 3.2. *The following assertions hold true:*

1) every solution of the (AI')-problem determines a solution of the (4B)-problem by the formula

$$\Omega_F h = (\Gamma + \Phi F_+) h, \quad h \in H; \quad (3.2.1)$$

2) conversely, an arbitrary solution of the (4B)-problem generates a solution of the (AI')-problem

$$F_{+, \Omega} h = \Phi^*(\Omega - \Gamma)h, \quad h \in H; \quad (3.2.2)$$

3) the mappings are bijective and mutually inverse.

Let us prove the first assertion of the proposition. In other words, we should demonstrate that Ω_F , defined by (3.2.1), possesses properties (2.2.2) and (2.2.3). We have

$$P_{H'} \Omega_F h = P_{H'} (\Gamma + \Phi F_+) h = P_{H'} \Gamma h = \Gamma h,$$

since $\Phi F_+ h \perp H'$. We proceed with the estimation of the norm of the operator Ω_F . We get

$$\begin{aligned} (\Omega_F h, \Omega_F h) &= ((\Gamma + \Phi F_+) h, (\Gamma + \Phi F_+) h) \\ &= (\Gamma h, \Gamma h) + (F_+ h, F_+ h). \end{aligned}$$

Here we have used the orthogonality of Γh to $\Phi F_+ h$, for every h from H , and the equality $\Phi^* \Phi = 1_F$. The latter sum can be estimated with the help of (3.1.2):

$$(\Gamma h, \Gamma h) + (F_+ h, F_+ h) \leq (h, h),$$

hence $\|\Omega_F\| \leq 1$. To show that the operator Ω_F is actually the multiplication by a certain operator-valued function, it suffices to verify that

$$\Omega_F t h = t \Omega_F h, \quad \forall h \in H.$$

Using commutativity relations (2.3.1) and (3.1.1), we obtain

$$\begin{aligned} (\Gamma + \Phi F_+) t h &= t \Gamma h - \Phi M_2 h + \Phi (t F_+ h + M_2 h) \\ &= t (\Gamma + \Phi F_+) h = t \Omega_F h. \end{aligned}$$

This computation completes the proof of the first part of the proposition.

Let us pass to the second part now. The formula for orthogonal projection $P_{H'}$ (see (2.3.4)) implies that

$$\Gamma h = \Omega h - \Phi P_+ \Phi^* \Omega h, \quad (3.2.3)$$

and, hence, by virtue of (3.2.2)

$$F_{+, \Omega} h = P_+ \Phi^* \Omega h \in H^2(F). \quad (3.2.4)$$

We obtain inequality (3.1.2) for $F_{+,\Omega}$ by estimating the norm of the functional l ,

$$l(g) = (g, F_{+,\Omega}h), \quad g \in L^2(F).$$

The Schwarz–Bunyakovskii inequality immediately yields that

$$|l(g)|^2 = |(\Phi g, (\Omega - \Gamma)h)|^2 \leq ((\Omega - \Gamma)h, (\Omega - \Gamma)h)(g, g).$$

It is not difficult to see that

$$((\Omega - \Gamma)h, (\Omega - \Gamma)h) = (\Omega h, \Omega h) - (\Gamma h, \Gamma h),$$

since $(\Omega h, \Gamma h) = (P_{H'}\Omega h, \Gamma h) = (\Gamma h, \Gamma h)$, in accordance with (2.2.2). Hence, we have

$$(\Omega h, \Omega h) - (\Gamma h, \Gamma h) \leq (h, h) - (\Gamma h, \Gamma h) = ((I - \Gamma^*\Gamma)h, h),$$

because of $\|\Omega\| \leq 1$. Thus we conclude that

$$\|l\|^2 = \|F_{+,\Omega}h\|^2 \leq ((I - \Gamma^*\Gamma)h, h), \quad \forall h \in H,$$

and the second statement is also proved.

Let Ω be an arbitrary solution of the (4B)-problem. We set $F_{+,\Omega}$ by Ω with the help of (3.2.2), or, which is the same, with the help of (3.2.4). Then, combining (3.2.1) and (3.2.3), we obtain

$$\tilde{\Omega}_{F_{+,\Omega}}h = \Gamma h + \Phi F_{+,\Omega}h = \Omega h - \Phi P_+ \Phi^* \Omega h + \Phi P_+ \Phi^* \Omega h = \Omega h, \quad h \in H.$$

Conversely, let a solution of the (AI')-problem, say, F_+ , be given. We construct the corresponding solution of the (4B)-problem according to (3.2.1). Then we pass one more time to a solution \tilde{F}_{Ω_F} of the (AI')-problem by (3.2.2). Summing up, we have

$$\tilde{F}_{\Omega_F} = \Phi^*(\Omega_F - \Gamma) = \Phi^*(\Gamma + \Phi F_+ - \Gamma) = F_+.$$

We are done.

3.3. First, being in the setting of the (AI')-problem, we are able to rewrite formula (1.2.2) in the terms of the (4B')-problem. Indeed, taking into account equalities (2.1.1), (1.2.2) and formula (2.3.2), we see that

$$d_V = \{tH^2(E_1) \oplus L^2(E_2)\} \oplus \{0\},$$

and

$$N_{d_V} = \{H^2(E_1) \oplus L^2(E_2)\} \ominus \{tH^2(E_1) \oplus L^2(E_2)\} = E_1 \ominus (E_1 \cap \text{Ker}(I - \Gamma^*\Gamma)^{1/2}).$$

Unfortunately, we cannot say anything about the structure of the second defect subspace.

Since all solutions of the (4B)-problem are, in fact, operator-valued functions from $L^2(E_1 \oplus E_2, E'_1 \oplus E'_2)$, it is sufficient to point out the values of the operators on vectors from $E_1 \oplus E_2$. Comparing the second row of (1.4.2) with (3.2.1), we conclude that the following parametrization takes place

$$\Omega_F \begin{bmatrix} e_1 \\ e_2 \end{bmatrix} = \left(\Gamma + \begin{bmatrix} \theta \\ 0 \end{bmatrix} G_{+2} + \begin{bmatrix} \theta \\ 0 \end{bmatrix} s_2(1_{\mathcal{N}_2} - \varepsilon s)^{-1} \varepsilon G_{+1} \right) \begin{bmatrix} e_1 \\ e_2 \end{bmatrix}, e_1 \in E_1, e_2 \in E_2,$$

where ε is an arbitrary function from $\mathbf{B}[\mathcal{N}_1, \mathcal{N}_2]$.

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**Абстрактная интерполяционная задача и задача
о четырех блоках**

С.А. Купин

Изучается 4-блок задача как абстрактная интерполяционная задача с подходящим образом подобранными данными. Этот подход позволяет получить описание всех решений произвольной 4-блок задачи.

**Абстрактна інтерполяційна задача та задача
про чотири блоки**

С.О. Купін

Вивчається 4-блок задача як абстрактна інтерполяційна задача з відповідно підібраними даними. Цей підхід дозволяє одержати опис усіх розв'язань довільної 4-блок задачі.